

MA612L-Partial Differential Equations

Lecture 21: Laplace Equation - Mean Value, Maximum-Minimum Principle and Liouville's Theorem

Panchatcharam Mariappan¹

¹Associate Professor
Department of Mathematics and Statistics
IIT Tirupati, Tirupati

September 29, 2025





Mean-Value Formulas

Mean-value Formulas



Now, we are looking again at the average integral, which we have discussed for the wave equation.

Theorem 1

$u \in C^2(\Omega)$ is harmonic, if and only if,

$$u(\mathbf{x}) = \oint_{\partial B(\mathbf{x}, r)} u dS = \oint_{B(\mathbf{x}, r)} u dy \quad (1)$$

for each ball $B(\mathbf{x}, r) \subset \Omega$.

Proof:



As we have seen in the Wave equation, we have

$$\phi(r) := \int_{\partial B(\mathbf{x}, r)} u(\mathbf{y}) dS(\mathbf{y}) = \int_{\partial B(\mathbf{0}, 1)} u(\mathbf{x} + r\mathbf{z}) dS(\mathbf{z})$$

$$\begin{aligned} \implies \phi'(r) &= \int_{\partial B(\mathbf{0}, 1)} Du(\mathbf{x} + r\mathbf{z}) \cdot \mathbf{z} dS(\mathbf{z}) \\ &= \int_{\partial B(\mathbf{x}, r)} Du(\mathbf{y}) \cdot \frac{\mathbf{y} - \mathbf{x}}{r} dS(\mathbf{y}) \\ &= \int_{\partial B(\mathbf{x}, r)} \frac{\partial u}{\partial \nu} dS(\mathbf{y}) \end{aligned}$$

Proof(continued)

By Green's formula, we have

$$\int_{\Omega} \Delta u d\mathbf{x} = \int_{\partial\Omega} \frac{\partial u}{\partial \nu} dS$$

Hence

$$\begin{aligned} \implies \phi'(r) &= \int_{\partial B(\mathbf{x},r)} \frac{\partial u}{\partial \nu} dS(\mathbf{y}) = \frac{1}{n\alpha(n)r^{n-1}} \int_{\partial B(\mathbf{x},r)} \frac{\partial u}{\partial \nu} dS(\mathbf{y}) \\ &= \frac{r}{n\alpha(n)r^n} \int_{B(\mathbf{x},r)} \Delta u d(\mathbf{y}) = \frac{r}{n} \int_{B(\mathbf{x},r)} \Delta u d(\mathbf{y}) = 0 \end{aligned}$$

$\implies \phi$ is constant.



Proof(continued)



$$\phi(r) = \lim_{t \rightarrow 0} \phi(t) = \lim_{t \rightarrow 0} \int_{\partial B(\mathbf{x}, t)} u(\mathbf{y}) dS(\mathbf{y}) = u(\mathbf{x})$$

Therefore, we have proved that

$$u(\mathbf{x}) = \int_{\partial B(\mathbf{x}, r)} u dS \implies \int_{\partial B(\mathbf{x}, r)} u dS = n\alpha(n)r^{n-1}u(\mathbf{x})$$

Now, it is sufficient to prove that

$$u(\mathbf{x}) = \int_{B(\mathbf{x}, r)} u dy$$

when u is harmonic

Proof(continued)



The following results can be obtained from the Coarea formula (or a curvilinear Fubini Theorem). For continuous integrable functions $f : \mathbb{R}^n \rightarrow \mathbb{R}$, the spherical integration formula is given by

$$\int_{\mathbb{R}^n} f d\mathbf{x} = \int_0^\infty \left(\int_{\partial B(\mathbf{x}, r)} f dS \right) dr$$

In particular,

$$\int_{B(\mathbf{x}, r)} f d\mathbf{x} = \int_0^r \left(\int_{\partial B(\mathbf{x}, s)} f dS \right) ds$$

"Integrating over a ball is nothing but integrating over each inner sphere starting from 0."

Proof(continued)



$$\begin{aligned}\oint_{B(\mathbf{x},r)} u dy &= \frac{1}{\alpha(n)r^n} \int_{B(\mathbf{x},r)} u dy = \frac{1}{\alpha(n)r^n} \int_0^r \left(\int_{\partial B(\mathbf{x},s)} u dS \right) ds \\ &= \frac{1}{\alpha(n)r^n} \int_0^r n\alpha(n)s^{n-1}u(\mathbf{x})ds = u(\mathbf{x})\end{aligned}$$

Conversely, suppose, u is not harmonic, then there exists a ball $B(\mathbf{x}, r) \subset \Omega$ such that $\Delta u \neq 0$. Suppose $\Delta u > 0$ within $B(\mathbf{x}, r)$. Then

$$0 = \phi'(r) = \frac{r}{n} \oint_{B(\mathbf{x},r)} \Delta u d(\mathbf{y}) > 0 \Rightarrow \Leftarrow$$

Hence the proof.



Maximum-Minimum Principle

Strong Maximum Principle



Theorem 2

Suppose $u \in C^2(\Omega) \cap C(\overline{\Omega})$ is harmonic within Ω .

1. If Ω is connected and there exists a point $x_0 \in \Omega$ such that

$$u(x_0) = \max_{\overline{\Omega}} u$$

then u is constant within Ω

2. Further

$$\max_{\overline{\Omega}} u = \max_{\partial\Omega} u$$

Proof:



Suppose there exists a point $x_0 \in \Omega$ with

$$u(x_0) = M := \max_{\overline{\Omega}} u$$

Then for $0 < r < \text{dist}(x_0, \partial\Omega)$, the mean value property asserts that

$$M = u(x_0) = \int_{B(\mathbf{x}_0, r)} u d\mathbf{y} \leq M$$

The equality holds only if $u \equiv M$ within $B(\mathbf{x}_0, r)$. Hence $u(\mathbf{y}) = M$ for all $\mathbf{y} \in B(\mathbf{x}, r)$. Therefore, $\Omega_1 = \{\mathbf{x} \in \Omega : u(\mathbf{x}) = M\}$ is open and relatively closed in Ω . Also $\Omega_1 = \Omega$ if Ω is connected. Hence, u is constant within Ω . The second part follows immediately.

Strong Maximum Principle:



Remarks

1. The first part of the theorem is the strong maximum principle
2. The second part of the theorem is called the maximum principle.
3. If we replace u by $-u$, we obtain strong minimum and minimum principles.
4. If $u = f$ on $\partial\Omega$ where $f \geq 0$, then u is positive everywhere in Ω if f is positive somewhere on $\partial\Omega$
5. We can prove the uniqueness of solutions of boundary value problems for the Poisson equation using this maximum principle.

Uniqueness



Theorem 3 (Uniqueness)

Let $f \in C(\partial\Omega)$, $g \in C(\Omega)$. Then there exists at most one solution $u \in C^2(\Omega) \cap C(\bar{\Omega})$ of the boundary value problem

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases} \quad (2)$$

Proof; Let u_1 and u_2 satisfy (2). Let $w = u_1 - u_2 \implies \Delta w = 0$. Now, we apply the strong maximum principle on this, and we obtain that $w \equiv 0$. Hence $u_1 = u_2$.

Thanks

Doubts and Suggestions

panch.m@iittp.ac.in

